

Semi-Lagrangian Methods for Advection-Diffusion of Differential Form

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joint work with R. Hiptmair (ETH Zürich)

Recent Advances on Theory and Applications
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Case Study: Magnetic Advection

Magnetoquasistatic electrodynamic equations in moving media:

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$$\begin{array}{ll} \text{Faraday's Law} & \mathbf{curl} \mathbf{E} = -\partial_t \mathbf{B}, \\ \text{Ampere's Law} & \mathbf{curl} \mathbf{H} = \mathbf{j}, \\ \text{Ohm's Law} & \mathbf{j} = \sigma(\mathbf{E} + \boldsymbol{\beta} \times \mathbf{B}), \\ & \mathbf{B} = \mu \mathbf{H}, \\ \text{Constraint} & \mathbf{div} \mathbf{B} = 0, \end{array}$$

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Second order operator and first order operator that is parametrized by velocity.

Scalar Advection-Diffusion Problem

Given: (continuous) velocity field $\beta : \Omega \mapsto \mathbb{R}^d$ on bounded domain $\Omega \subset \mathbb{R}^d$

$$\beta \cdot \mathbf{grad} u - \operatorname{div}(\varepsilon \mathbf{grad} u) = f \quad \text{in } \Omega .$$

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$$k\text{-form } \omega \in \Lambda^k(\Omega) \hat{=} \text{ mapping } \omega : \left\{ \begin{array}{l} \text{Oriented} \\ k\text{-dimensional} \\ \text{sub-manifolds } \subset \Omega \end{array} \right. \mapsto \mathbb{R}$$

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In 3D: equivalent **vector proxy formulation**

$$\begin{array}{ll}
 k = 0 & : \quad -\operatorname{div}(\varepsilon \mathbf{grad} u) = f \quad \rightarrow \text{diffusion} \\
 k = 1 & : \quad \mathbf{curl}(\varepsilon \mathbf{curl} u) = \mathbf{f} \quad \rightarrow \text{magnetostatics} \\
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advection term \rightarrow $\beta \cdot \mathbf{grad} u$
 diffusive term \rightarrow $\operatorname{div}(\varepsilon \mathbf{grad} u)$
 Sobolev space \rightarrow u

Recall: general 2nd-order boundary value problem in $H\Lambda^k(d, \Omega)$

$$\delta(\varepsilon d\omega) = \varphi \quad \text{in } \Omega.$$

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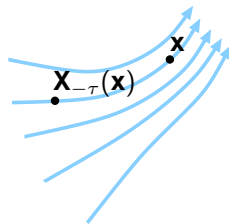
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Lie Derivatives

Let $\mathbf{X}_t : \Omega \rightarrow \Omega \hat{=}$ flow induced by $\beta = \beta(\mathbf{x})$:



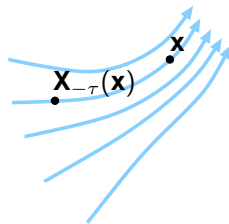
Special case $d = 3$, Lie derivatives:

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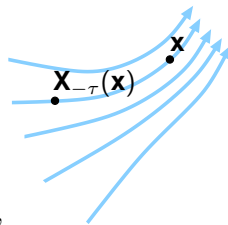
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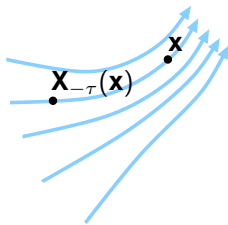
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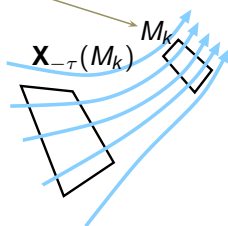
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$$\begin{aligned} \int_{M_k} \mathbf{L}_\beta \omega &:= \lim_{\tau \rightarrow 0} \frac{1}{\tau} \left(\int_{M_k} \omega - \int_{\mathbf{X}_{-\tau}(M_k)} \omega \right) \\ &= \lim_{\tau \rightarrow 0} \frac{1}{\tau} \left(\int_{M_k} \omega - \int_{M_k} \mathbf{X}_{-\tau}^* \omega \right) \\ &= \int_{M_k} \frac{d}{dt} \mathbf{X}_{-\tau}^* \omega \end{aligned}$$

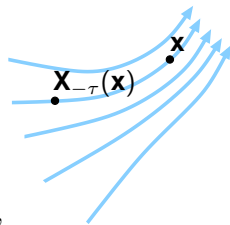
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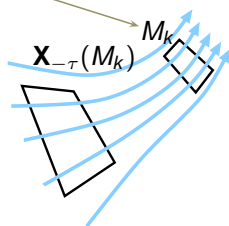
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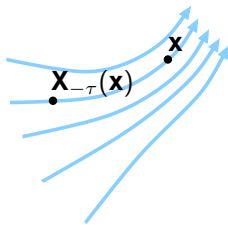
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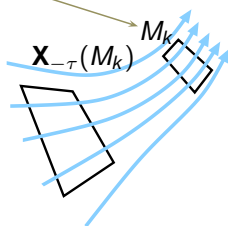
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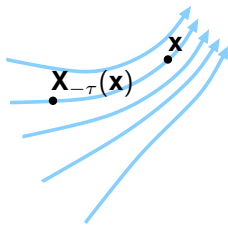
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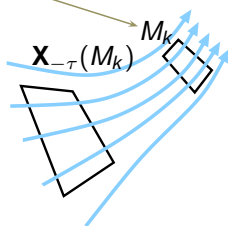
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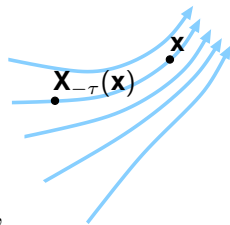
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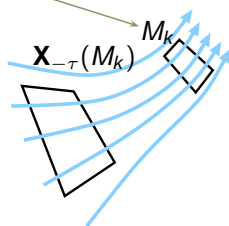
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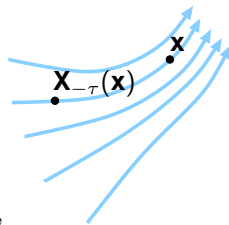


Cartan's formula: $\mathbf{L}_\beta \omega := \frac{d}{dt} \mathbf{X}_{-\tau}^* \omega = \mathbf{d} i_\beta \omega + i_\beta \mathbf{d} \omega.$

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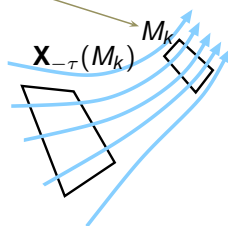
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$$\begin{aligned} \int_{M_k} \mathbf{L}_\beta \omega &:= \lim_{\tau \rightarrow 0} \frac{1}{\tau} \left(\int_{M_k} \omega - \int_{\mathbf{x}_{-\tau}(M_k)} \omega \right) \\ &= \lim_{\tau \rightarrow 0} \frac{1}{\tau} \left(\int_{M_k} \omega - \int_{M_k} \mathbf{X}_{-\tau}^* \omega \right) \\ &= \int_{M_k} \frac{d}{dt} \mathbf{X}_{-\tau}^* \omega \end{aligned}$$

contraction

Cartan's formula: $\mathbf{L}_\beta \omega := \frac{d}{dt} \mathbf{X}_{-\tau}^* \omega = \mathbf{d} i_\beta \omega + i_\beta \mathbf{d} \omega.$



Generalized Advection-Diffusion Problem

Find $\omega = \omega(t) \in \Lambda^k(\Omega)$, $\Omega \subset \mathbb{R}^d$, such that

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MagnetoHydroDynamics
with constraint $\text{div} \mathbf{u} = 0$

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= **finite elements** for differential forms: $\Lambda_h^k(\mathcal{T}) \subset H\Lambda^k(\mathbf{d}, \Omega)$

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simplicial mesh

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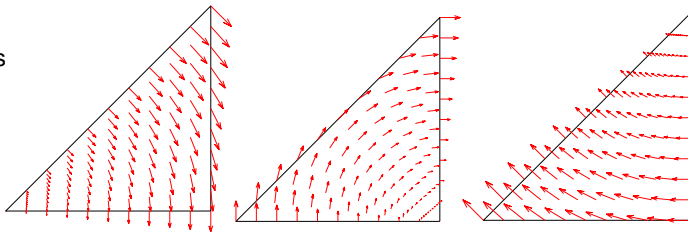
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vector proxies
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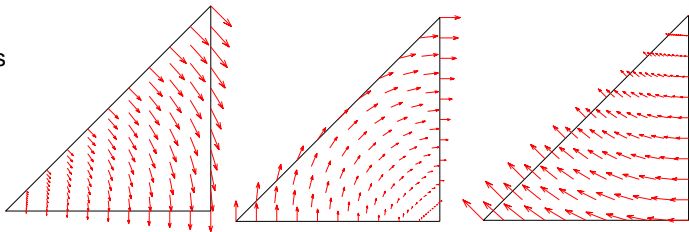
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Theorem: ($d > k > 0$) For $r \geq 0$ there are $\Lambda_h^k(\mathcal{T}) \subset H\Lambda^k(d, \Omega)$ that contain only piecewise polynomials of degree $s \leq r$ and:

$$\|\omega - I^k \omega\|_{L^2 \Lambda^k} = O(h^{r+1}) \quad \text{and} \quad \|d\omega - I^{k+1} d\omega\|_{L^2 \Lambda^{k+1}} = O(h^{r+1}).$$

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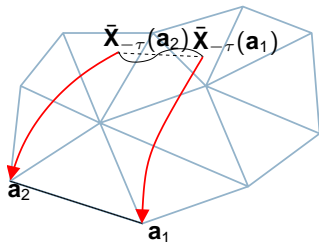
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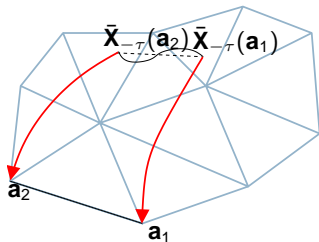
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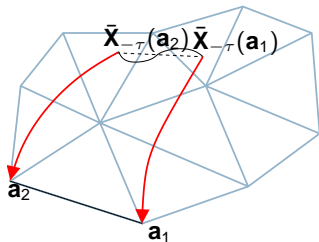
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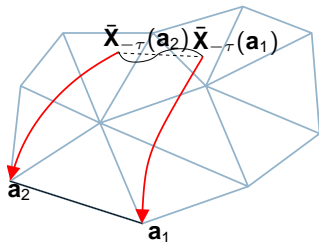
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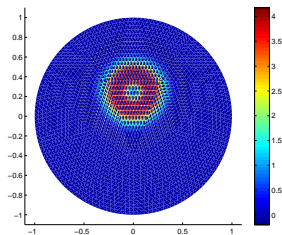
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Numerical Experiment

$$\partial_t \mathbf{u} + \beta \operatorname{div} \mathbf{u} + \operatorname{curl}(\mathbf{u} \times \beta) = 0 \quad \text{in unitcircle,}$$
$$\mathbf{u}(0) = (\text{localized bump})^2,$$

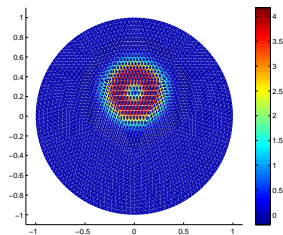
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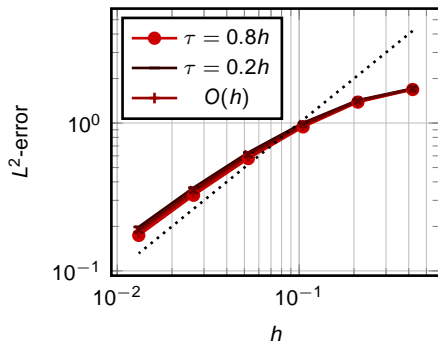
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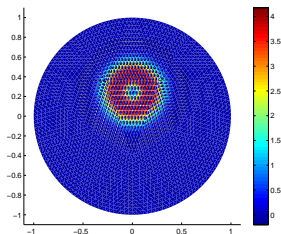
Convergence



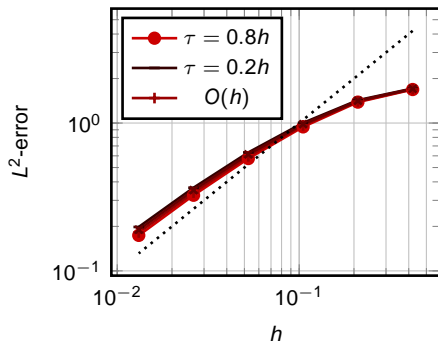
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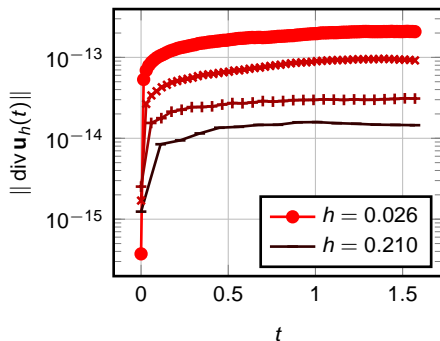
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Convergence



Preservation of closedness



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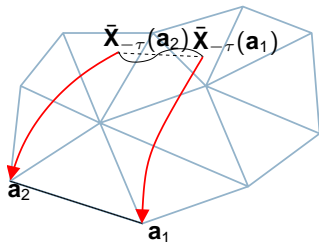
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Our Favourite SL-Scheme

Focus: $\Lambda_h^k(\mathcal{T}) \subset H\Lambda^k(d, \Omega)$ lowest order $r = 0$

Interpolation-based semi-Lagrangian scheme:

$$\omega_h^n = I^k \mathbf{X}_{-\tau}^* \omega_h^{n-1} + \tau I^k \varphi(t^n) \quad \text{for} \quad \partial_t \omega + \mathbf{L}_\beta \omega = \varphi.$$

canonical interpolation operator

Canonical degrees of freedom (d.o.f.) of $\Lambda_h^k(\mathcal{T})$:

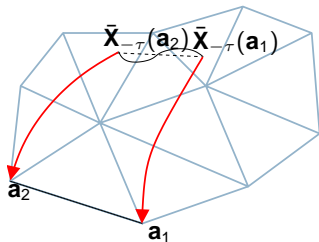
$$\text{set of dofs} = \left\{ \int_s \cdot : \Lambda^k(\Omega) \mapsto \mathbb{R}, s \text{ is subsimplex of } \mathcal{T} \right\}.$$

Advantages:

- ▶ preserves $d\omega_h^0$, i.e. if $\phi = 0$:
 $d\omega_h^n = d I^k \mathbf{X}_{-\tau}^* \omega_h^{n-1} = I^k \mathbf{X}_{-\tau}^* d\omega_h^{n-1}$;
- ▶ converges;

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$$\int_s \mathbf{X}_{-\tau}^* \omega = \int_{\mathbf{X}_{-\tau}(s)} \omega$$



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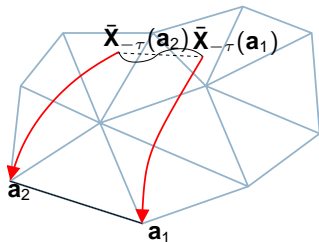
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Disadvantage:

- ▶ lack of L^2/L^∞ -stability property,
 $\|I^k \mathbf{X}_{-\tau}^* \omega_h\| \leq (1 + C\tau) \|\omega_h\| \quad k > 0$,
 hence no proof of convergence!

$$\int_s \mathbf{X}_{-\tau}^* \omega = \int_{\mathbf{X}_{-\tau}(s)} \omega$$



Advection-Diffusion: Convergence Theory

Find $\omega = \omega(t) \in \Lambda^k(\Omega)$, $\Omega \subset \mathbb{R}^d$, such that

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Can we use a similar argument for SL-methods?

- ▶ yes, at least for SL-Methods based on L^2 -projection.

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Non-stationary advection: Find $\omega = \omega(t) \in \Lambda^k(\Omega)$, $\Omega \subset \mathbb{R}^d$, such that

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Assumptions: $\left\{ \begin{array}{l} \mathbf{P}_h \mathbf{X}_{-\tau}^* \omega_h \text{ can be computed exactly,} \end{array} \right.$

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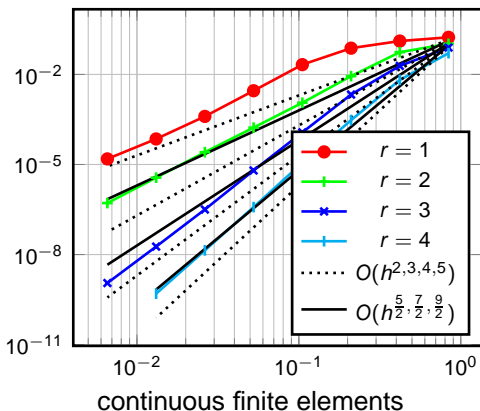
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rotating bump on unit-circle, $\beta = (-y, x)$, smooth initial data, $\tau = \frac{0.8}{\sqrt{2}}h$

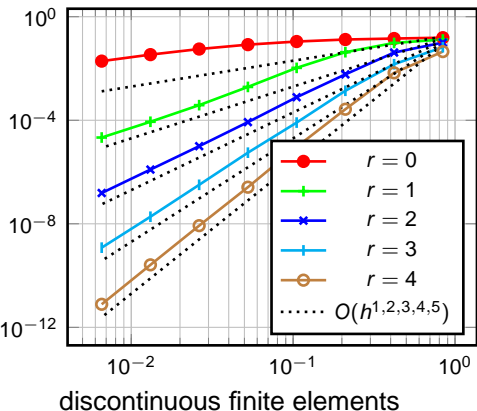
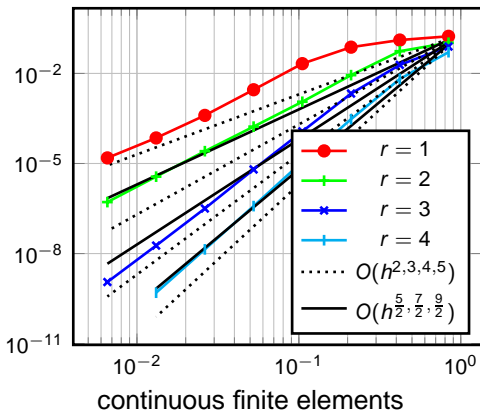
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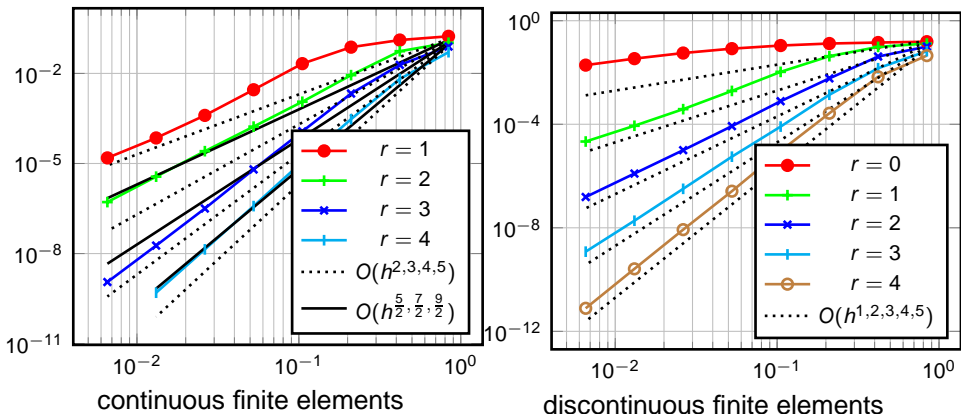
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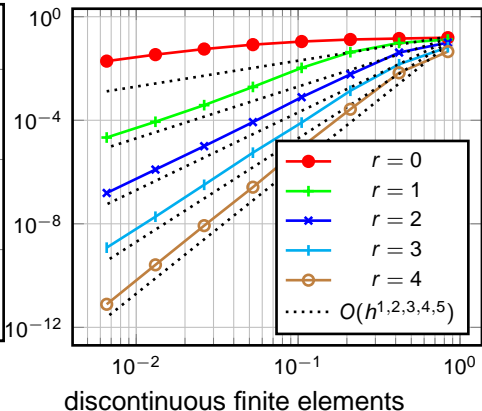
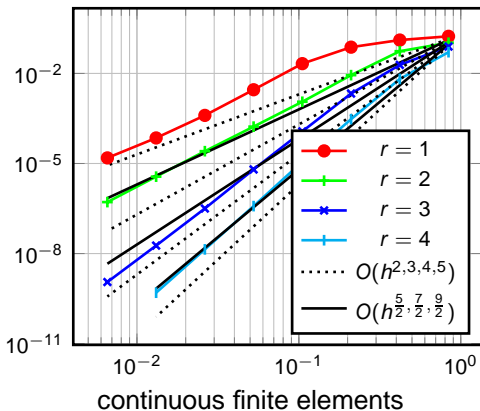
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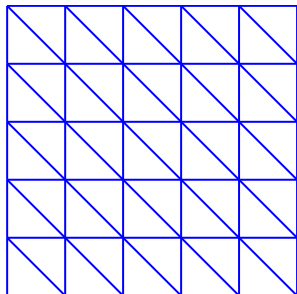
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Computation (approximation) of $P_h \mathbf{X}_{-\tau}^* \omega_h^{n-1}$ is the essence of SL-schemes!

Fully Discrete SL-Methods

Approximation of $P_h \mathbf{X}_{-\tau}^* \omega_h^{n-1}$:

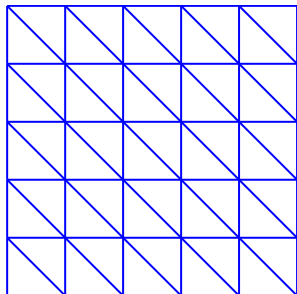
$$(\omega_h^n, \eta_h)_\Omega = \left(\mathbf{X}_{-\tau}^* \omega_h^{n-1}, \eta_h \right)_\Omega + \dots, \quad \forall \eta_h \in \Lambda_h^k(\mathcal{T})$$



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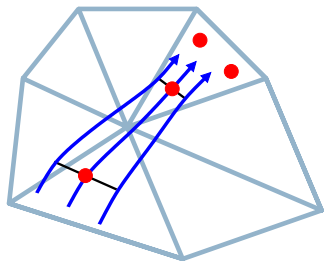


► FEM-Quadrature on \mathcal{T} :

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Numerical Experiment: Pure advection ($k = 1$, $d = 2$) and quadrature

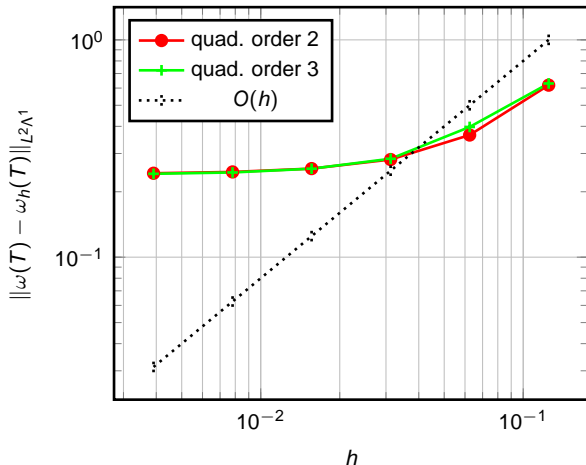
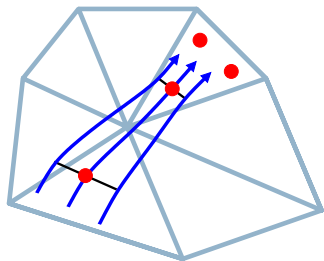
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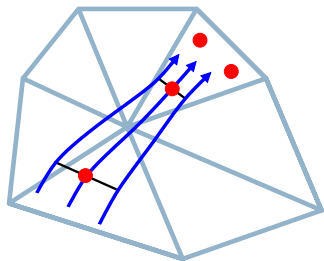
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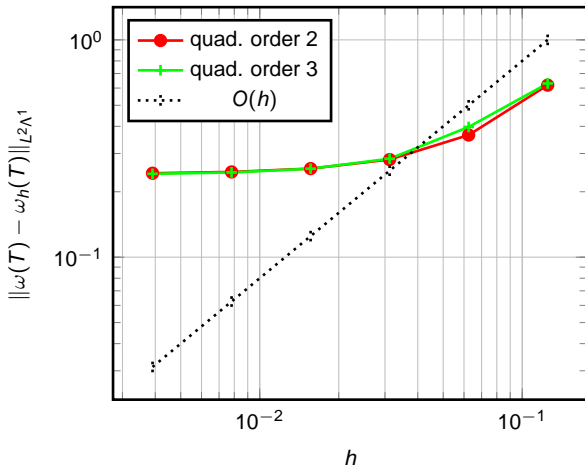
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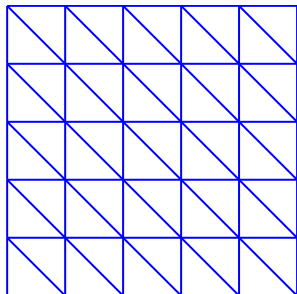
No convergence, why?



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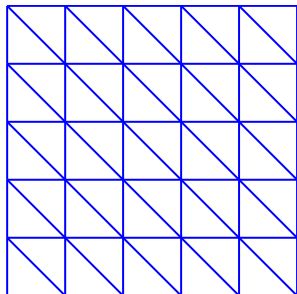


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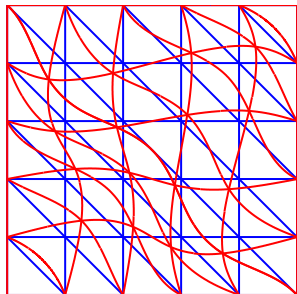


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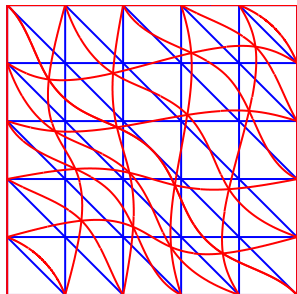


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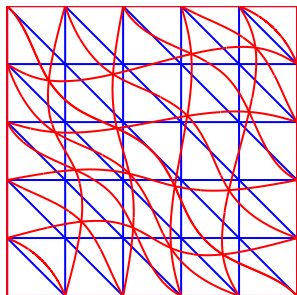


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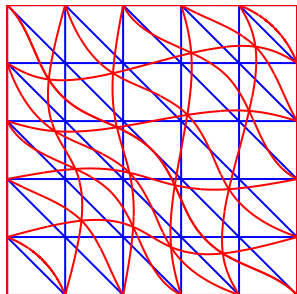
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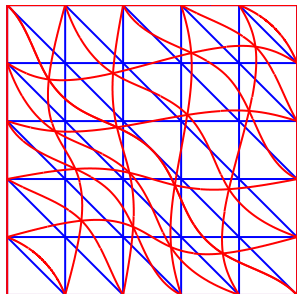
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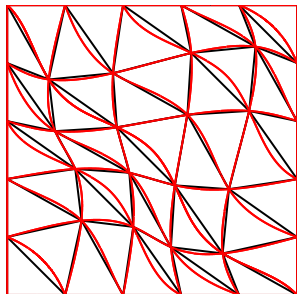
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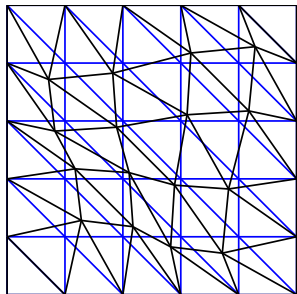
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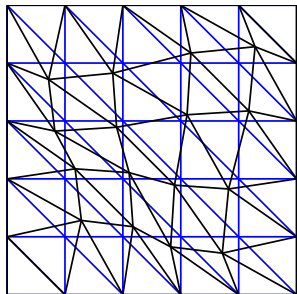
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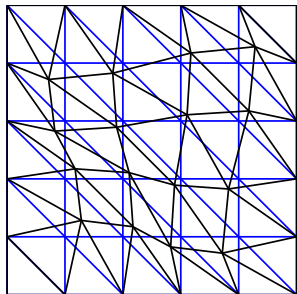
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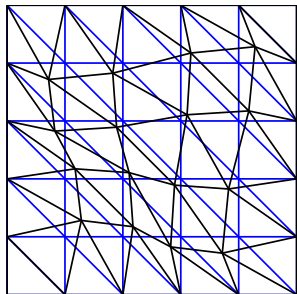
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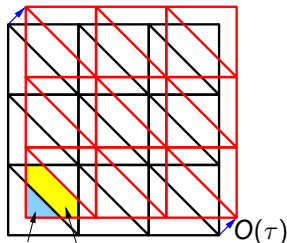
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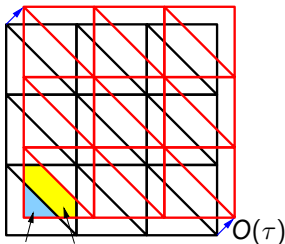
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For small τ Semi-Lagrange is close to explicit Euler, and $\text{error} = O(\tau + h^{r+\frac{1}{2}})$ for explicit Euler!

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characteristic methods for stationary advection:

stability: $\mathbf{C}_1 \|\eta_h\|_{h,\tau}^2 \leq \mathbf{a}_\tau(\eta_h, \eta_h),$

continuity: $\mathbf{a}_\tau(\omega, \eta_h) \leq \mathbf{C}_2 \tau^{-\frac{1}{2}} \|\omega\|_{L^2\Lambda^k} \|\eta_h\|_{h,\tau},$

consistency: $|\mathbf{a}_\tau(\omega, \eta_h) - (\mathbf{L}_\beta \omega, \eta_h)_\Omega| \leq \mathbf{C}_3 \tau \|\omega\|_{H^2\Lambda^k} \|\eta_h\|_{L^2\Lambda^k},$

convergence: $\|\omega - \omega_h\|_{h,\tau} \leq \mathbf{C} \left(h^{r+1} \tau^{-\frac{1}{2}} + \tau \right) \|\omega\|_{H^{r+1}\Lambda^k(\Omega)},$

if $\frac{1}{2\tau} (\omega, \omega)_\Omega - \frac{1}{2\tau} (\mathbf{X}_{-\tau}^* \omega, \mathbf{X}_{-\tau}^* \omega)_\Omega$ positive .

Advection-Diffusion II

Main tool is analysis of characteristic methods for stationary advection:

$$\mathbf{a}_\tau(\omega_h, \eta_h) := \frac{1}{\tau} (\omega_h, \eta_h)_\Omega - \frac{1}{\tau} (\mathbf{X}_{-\tau}^* \omega_h, \eta_h)_\Omega = l(\eta_h), \quad \tau \text{ artificial parameter}$$

in some mesh and τ -dependent norm:

$$\|\omega\|_{h,\tau}^2 := \|\omega\|_{L^2\Lambda^k}^2 + \frac{1}{2\tau} \|\omega - \mathbf{X}_{-\tau}^* \omega\|_{L^2\Lambda^k}^2 \xrightarrow{\tau \rightarrow 0} \|\omega\|_{\text{DG}}^2 \geq \|\omega\|_{L^2\Lambda^k}^2.$$

characteristic methods for stationary advection:

e.g. $k = 0$:
 $\text{div } \beta > 0$

stability: $C_1 \|\eta_h\|_{h,\tau}^2 \leq \mathbf{a}_\tau(\eta_h, \eta_h),$

continuity: $\mathbf{a}_\tau(\omega, \eta_h) \leq C_2 \tau^{-\frac{1}{2}} \|\omega\|_{L^2\Lambda^k} \|\eta_h\|_{h,\tau},$

consistency: $|\mathbf{a}_\tau(\omega, \eta_h) - (\mathbf{L}_\beta \omega, \eta_h)_\Omega| \leq C_3 \tau \|\omega\|_{H^2\Lambda^k} \|\eta_h\|_{L^2\Lambda^k},$

convergence: $\|\omega - \omega_h\|_{h,\tau} \leq C \left(h^{r+1} \tau^{-\frac{1}{2}} + \tau \right) \|\omega\|_{H^{r+1}\Lambda^k(\Omega)},$

if $\frac{1}{2\tau} (\omega, \omega)_\Omega - \frac{1}{2\tau} (\mathbf{X}_{-\tau}^* \omega, \mathbf{X}_{-\tau}^* \omega)_\Omega$ positive $\xrightarrow{\tau \rightarrow 0}$ $\mathbf{L}_\beta + \mathcal{L}_\beta$ positive.

Advection-Diffusion II

Main tool is analysis of **characteristic methods** for **stationary advection**:

$$a_\tau(\omega_h, \eta_h) := \frac{1}{\tau} (\omega_h, \eta_h)_\Omega - \frac{1}{\tau} (\mathbf{X}_{-\tau}^* \omega_h, \eta_h)_\Omega = l(\eta_h), \quad \tau \text{ artificial parameter}$$

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characteristic methods for stationary advection:

e.g. $k = 0$:
 $\text{div } \beta > 0$

stability: $C_1 \|\eta_h\|_{h,\tau}^2 \leq a_\tau(\eta_h, \eta_h),$

continuity: $a_\tau(\omega, \eta_h) \leq C_2 \tau^{-\frac{1}{2}} \|\omega\|_{L^2\Lambda^k} \|\eta_h\|_{h,\tau},$

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convergence: $\|\omega - \omega_h\|_{h,\tau} \leq C \left(h^{r+1} \tau^{-\frac{1}{2}} + \tau \right) \|\omega\|_{H^{r+1}\Lambda^k(\Omega)},$

if $\frac{1}{2\tau} (\omega, \omega)_\Omega - \frac{1}{2\tau} (\mathbf{X}_{-\tau}^* \omega, \mathbf{X}_{-\tau}^* \omega)_\Omega$ positive $\xrightarrow{\tau \rightarrow 0}$ $L_\beta + \mathcal{L}_\beta$ positive.

Ritz-Galerkin projector with respect to **characteristic methods** yields,

SL for advection-diffusion of k -forms $k > 0$: $\text{error} = \left(h^{r+1} \tau^{-\frac{1}{2}} + \tau \right)$

Summary and Conclusions

Summary:

- ▶ Lagrangian timestepping schemes for differential forms;
- ▶ Analysis of fully discrete SL-schemes;
- ▶ Proof of uniform convergence for advection-diffusion;

Outlook:

- ▶ Analysis of interpolation based Semi-Lagrangian schemes;
- ▶ Constraint preserving schemes for hyperbolic problems;

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