

Fourier Jacobi-Bessel series for Meijer's G-function

S.D. BAJPAI

RIASSUNTO – Si introduce una nuova classe delle serie di Fourier Jacobi-Bessel per la G-funzione di Meijer.

ABSTRACT – In this paper, we present a new class of Fourier Jacobi-Bessel series for Meijer's G-function.

KEY WORDS – Fourier Jacobi-Bessel series - Meijer's G-function - Two-dimensional partial differential equation.

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1 – Introduction

The object of this paper is to introduce a new class of Fourier Jacobi-Bessel series for Meijer's G-function [3, pp. 206-222], and present few of its particular cases. We also show that our Fourier Jacobi-Bessel Series is related to the solution of a two dimensional partial differential equation.

The following formulae are required in the proof:

The integral [1, p. 177, (2.1)]:

$$(1.1) \quad \int_{-1}^1 (1-x)^{\rho}(1+x)^{\beta} P_u^{(\alpha, \beta)}(x) G_{p,q}^{m,n} \left[z(1-x)^{\delta} \left| \begin{matrix} a_p \\ b_q \end{matrix} \right. \right] dx$$

$$= \frac{2^{\beta+\rho+1} \Gamma(\beta+u+1)}{\delta^{\beta+1} u!} G_{p+2\delta, q+2\delta}^{m+\delta, n+\delta} \left[z 2^{\delta} \left| \begin{matrix} \Delta(\delta, -\rho), a_p, \Delta(\delta, \alpha-\rho) \\ \Delta(\delta, \alpha-\rho+u), b_q, \Delta(\delta, -1-\beta-\rho-u) \end{matrix} \right. \right],$$

where δ is a positive integer, $2(m+n) > p+q$, $|\arg z| < (m+n-\frac{1}{2}p-\frac{1}{2}q)\pi$, $\operatorname{Re} \beta > -1$, $\operatorname{Re}(\rho + \delta b_j) > -1$ ($j = 1, 2, \dots, m$).

The integral [2, p. 37, (2.1)]:

$$(1.2) \quad \int_0^\infty y^{\sigma-1} J_\nu(y) Y_\mu(y) G_{p,q}^{m,n} \left[zy^{2\lambda} \left| \begin{matrix} a_p \\ b_q \end{matrix} \right. \right] = \frac{2^{\sigma-1} \lambda^{2(\sigma-1)}}{(2\pi)^{1-\lambda}} \cdot G_{p+5\lambda, q+\lambda}^{m, n+2\lambda} \left[2^{2\lambda} \lambda^{4\lambda} z \left| \begin{matrix} \Delta(\lambda, \frac{2-\sigma-\mu-\nu}{2}), \Delta(\lambda, \frac{2-\sigma+\mu-\nu}{2}), a_p, \\ \Delta(\lambda, \frac{3-\nu+\mu-\sigma}{2}), \Delta(\lambda, \frac{\nu-\mu-\sigma+2}{2}), \Delta(\lambda, \frac{\nu+\mu-\sigma+2}{2}), \\ b_q, \Delta(\lambda, \frac{3-\nu+\mu-\sigma}{2}) \end{matrix} \right. \right],$$

where λ is a positive integer, $2(m+n) > p+q$, $|\arg z| < (m+n-\frac{1}{2}p-\frac{1}{2}q)\pi$, $\operatorname{Re}(\sigma + \nu \pm \mu + 2\lambda b_j) > 0$ ($j = 1, 2, \dots, m$), $\operatorname{Re}(\sigma + 2\lambda(a_j - 1)) < 1$ ($j = 1, 2, \dots, n$).

The orthogonality property of the Jacobi polynomials [4, p. 285, (5) and (9)]:

$$(1.3) \quad \int_{-1}^1 (1-x)^\alpha (1+x)^\beta P_n^{(\alpha, \beta)}(x) P_m^{(\alpha, \beta)}(x) dx = \begin{cases} 0 & \text{if } m \neq n, \\ \frac{2^{\alpha+\beta+1} \Gamma(\alpha+n+1) \Gamma(\beta+n+1)}{n! (\alpha+\beta+2n+1) \Gamma(\alpha+\beta+n+1)} & \text{if } m = n; \end{cases}$$

where $\operatorname{Re} \alpha > -1$, $\operatorname{Re} \beta > -1$.

The orthogonality property of the Bessel functions [5, p. 291, (6)]:

$$(1.4) \quad \int_0^\infty x^{-1} J_{\nu+2n+1}(x) J_{\nu+2m+1}(x) dx = \begin{cases} 0 & \text{if } m \neq n, \\ (4n+2\nu+2)^{-1} & \text{if } m = n, \quad \operatorname{Re} \nu + m + 1 > -1. \end{cases}$$

2 - Fourier Jacobi-Bessel series

The Fourier Jacobi-Bessel series to be established is

$$\begin{aligned}
 (2.1) \quad & (1-x)^{\rho} y^{\sigma} Y_{\mu}(y) G_{p,q}^{m,n} \left[z(1-x)^{\delta} y^{2\lambda} \left| \begin{matrix} a_p \\ b_q \end{matrix} \right. \right] \\
 &= \frac{2^{\rho+\sigma} \lambda^{2(\sigma-1)}}{\delta^{\beta+1} (2\pi)^{1-\lambda}} \sum_{\substack{r=0 \\ t=0}}^{\infty} \frac{\alpha + \beta + 2r + 1}{\Gamma(\alpha + r + 1)} \frac{\Gamma(\alpha + \beta + r + 1)(\nu + 2t + 1)}{\Gamma(\alpha + r + 1)} \\
 &\cdot P_r^{(\alpha, \beta)}(x) J_{\nu+2t+1}(y) \cdot \\
 &\cdot G_{p+2\delta+5\lambda, q+2\delta+\lambda}^{m+\delta, n+\delta+2\lambda} \left[\frac{z^{2\delta}}{(2\lambda^2)^{-2\lambda}} \left[\begin{matrix} \Delta(\delta, -\alpha - \rho), \Delta(\lambda, \frac{1-\sigma-\mu-\nu-2t}{2}), \\ \Delta(\lambda, \frac{1-\sigma+\mu-\nu-2t}{2}), a_p, \Delta(\delta, -\rho), \\ \Delta(\lambda, \frac{2-\nu-2t+\mu-\sigma}{2}), \\ \Delta(\lambda, \frac{\nu+2t-\mu-\sigma+3}{2}), \Delta(\lambda, \frac{\nu+2t+\mu-\sigma+3}{2}), \\ \Delta(\delta, r-\rho), b_q, \Delta(\lambda, \frac{2-\nu-2t+\mu-\sigma}{2}), \Delta(\delta, -1-\alpha-\beta-\rho-r) \end{matrix} \right] \right]
 \end{aligned}$$

valid under the conditions of (1.1), (1.2), (1.3) and (1.4).

PROOF. Let

$$\begin{aligned}
 (2.2) \quad & f(x, y) = (1-x)^{\rho} y^{\sigma} Y_{\mu}(y) G_{p,q}^{m,n} \left[z(1-x)^{\delta} y^{2\lambda} \left| \begin{matrix} a_p \\ b_q \end{matrix} \right. \right] \\
 &= \sum_{\substack{r=0 \\ t=0}}^{\infty} C_{r,t} P_r^{(\alpha, \beta)}(x) J_{\nu+2t+1}(y).
 \end{aligned}$$

Equation (2.2) is valid, since $f(x, y)$ is continuous and of bounded variation in the region $-1 < x < 1$, $0 < y < \infty$. Multiplying both sides of (2.2) by $y^{-1} J_{\nu+2v+1}(y)$, integrating with respect to y from 0 to ∞ , and using (1.2) and (1.4), then multiplying both sides of the resulting expression by $(1-x)^{\alpha} (1+x)^{\beta} P_u^{(\alpha, \beta)}(x)$, integrating with respect to x from -1 to 1 and using (1.1) and (1.3), the value of $C_{r,t}$ is obtained. Substituting this value of $C_{r,t}$ in (2.2), the Fourier Jacobi-Bessel series

(2.1) is established.

NOTE: On applying the same procedure as above, we can establish three other forms of two-dimensional expansions of this class with the help of alternative forms of (1.1) and (1.2).

3 – Particular cases

Since on specializing the parameters Meijer's G-function yields almost all special functions appearing in applied mathematics and physical sciences [3]. Therefore, the result presented in this paper is of a general character and hence may encompass several cases of interest. However, we present below only few particular cases of our Fourier Jacobi-Bessel series.

In (2.1), putting $\delta = \lambda = 1$ and using the identity [1, p. 180, (2.6)], it reduces to the form

$$\begin{aligned}
 (3.1) \quad & (1-x)^{\rho} y^{\sigma} Y_{\mu}(y) G_{p,q}^{m,n} \left[z(1-x)y^2 \left| \begin{matrix} a_p \\ b_q \end{matrix} \right. \right] \\
 &= 2^{\rho+\sigma} \sum_{\substack{r=0 \\ t=0}}^{\infty} \frac{(-1)^r (\alpha + \beta + 2r + 1) \Gamma(\alpha + \beta + r + 1) (\nu + 2t + 1)}{\Gamma(\alpha + r + 1)} \\
 &\quad \cdot P_r^{(\alpha+\beta)}(x) J_{\nu+2t+1}(y) \cdot \\
 &\quad \cdot G_{p+7,q+3}^{m,n+4} \left[8z \left| \begin{matrix} -\rho, -\alpha-\rho, \frac{1-\alpha-\mu-\nu-2t}{2}, \frac{1-\sigma+\mu-\nu-2t}{2}, a_p, \\ \frac{2-\nu-2t+\mu-\sigma}{2}, \frac{\nu+2t-\mu-\sigma+3}{2}, \frac{\nu+2t+\mu-\sigma+3}{2}, \\ b_q, r-\rho, -1-\alpha-\beta-\rho, 2-\nu-2t+\mu-\sigma \end{matrix} \right. \right],
 \end{aligned}$$

valid under the conditions analogous to (2.1).

In (3.1), setting $m = 1$, $n = p$, $b_1 = 0$ and using [4, p. 439, (4)], we get

$$\begin{aligned}
 (3.2) \quad & (1-x)^{\rho} y^{\sigma} Y_{\mu}(y) {}_pF_{q-1} \left[\begin{matrix} 1-a_1, \dots, 1-a_p; -z(1-x) = y^2 \\ 1-b_2, \dots, 1-b_q; \end{matrix} \right] \\
 &= 2^{\rho+\sigma} \sum_{\substack{r=0 \\ t=0}}^{\infty} \frac{(-1)^r (\alpha + \beta + 2r + 1) \Gamma(\alpha + \beta + r + 1) (\nu + 2t + 1)}{\Gamma(\alpha + r + 1)} \cdot \\
 &\quad \cdot P_r^{(\alpha+\beta)}(x) J_{\nu+2t+1}(y) \cdot \\
 &\quad \cdot \frac{\Gamma(1+\rho) \Gamma(1+\alpha+\rho) \Gamma(\frac{1+\sigma+\mu+\nu+2t}{2}) \Gamma(\frac{1+\sigma-\mu+\nu+2t}{2})}{\Gamma(1-r+\rho) \Gamma(2+\alpha+\beta+\rho+r) \Gamma(\nu+2t+\sigma-\mu-1) \Gamma(\frac{2-\nu-2t+\mu-\sigma}{2})} \cdot \\
 &\quad \cdot \frac{1}{\Gamma(\frac{\nu+2t-\mu-\sigma+3}{2}) \Gamma(\frac{\nu+2t+\mu-\sigma+3}{2})} \cdot \\
 &\quad \cdot {}_{p+7}F_{q+2} \left[\begin{matrix} 1-a_1, \dots, 1-a_p, 1+\rho, 1+\alpha+\rho, \frac{1+\sigma+\mu+\nu+2t}{2}, \\ \frac{1+\sigma-\mu+\nu+2t}{2}, \frac{1+\nu+2t-\mu+\sigma}{2}, \frac{\mu+\sigma-\nu-2t-1}{2}, \\ \frac{\sigma-\nu-\nu-2t-1}{2}; \\ 1-b_2, \dots, 1-b_q, 1-r+\rho, 2+\alpha+\beta+\rho+r, \nu+2t+\sigma-r-1; -8x \end{matrix} \right],
 \end{aligned}$$

valid under the conditions analogous to (2.1).

In (3.2), putting $z = 0$, we obtain

$$\begin{aligned}
 (3.3) \quad & (1-x)^{\rho} y^{\sigma} Y_{\mu}(y) = 2^{\rho+\sigma} \sum_{\substack{r=0 \\ t=0}}^{\infty} \frac{(-1)^r (\alpha + \beta + 2r + 1)}{\Gamma(\alpha + r + 1)} \cdot \\
 &\quad \cdot \frac{\Gamma(\alpha + \beta + r + 1) (\nu + 2t + 1) \Gamma(1+\rho) \Gamma(1+\alpha+\rho) \Gamma(\frac{1+\sigma+\mu+\nu+2t}{2})}{\Gamma(1-r+\rho) \Gamma(2+\alpha+\beta+\rho+r) \Gamma(\nu+2t+\sigma-\mu-1)} \cdot \\
 &\quad \cdot \frac{\Gamma(\frac{1+\sigma-\mu+\nu+2t}{2})}{\Gamma(\frac{2-\nu-2t+\mu-\sigma}{2})} P_r^{(\alpha,\beta)}(x) J_{\nu+2t+1}(y),
 \end{aligned}$$

valid under the conditions analogous to (2.1).

4 – Two-dimensional partial differential equation

Let us consider the following two-dimensional partial differential equation

$$(4.1) \quad \frac{\partial u}{\partial t} = (1-x^2) \frac{\partial^2 u}{\partial x^2} + [\beta - \alpha - (\alpha + \beta + 2)x] \frac{\partial u}{\partial x} + y^2 \frac{\partial^2 u}{\partial y^2} + y \frac{\partial u}{\partial y} + y^2 u,$$

where $u \equiv u(x, y, t)$, and $u(x, y, 0) = f(x)$.

To solve (4.1), let us assume that (4.1) has a solution of the form:

$$(4.2) \quad u(x, y, t) = e^{-[n(n+\alpha+\beta+1)-m^2]t} X(x)Y(y).$$

The substitution of (4.2) into (4.1) yields the differential equation:

$$(4.3) \quad Y(y)[(1-x^2)X'' + \{\beta - \alpha - (\alpha + \beta + 2)x\}X' + n(n + \alpha + \beta + 1)X] + X(x)[y^2Y'' + yY' + (y^2 - m^2)Y] = 0.$$

We have, Jacobi equation

$$(4.4) \quad (1-x^2)X'' + \{\beta - \alpha - (\alpha + \beta + 2)x\}X' + n(n + \alpha + \beta + 1)X = 0,$$

with its solution $X = P_n^{(\alpha, \beta)}(x)$.

And, we have, Bessel equation

$$(4.5) \quad y^2Y'' + yY' + (y^2 - m^2)Y = 0,$$

with its solution $Y = J_m(y)$.

In view of (4.4) and (4.5), we conclude that to each eigenvalue given by (4.3), there corresponds the solution of (4.1), called an eigenfunction or eigenstate given by

$$(4.6) \quad u(x, y, t) = e^{-[n(n+\alpha+\beta+1)-m^2]t} P_n^{(\alpha, \beta)}(x) J_m(y).$$

In view of the principle of superposition, the solution of (4.1) is given by

$$(4.7) \quad u(x, y, t) = \sum_{\substack{n=0 \\ m=0}}^{\infty} A_{n,m} e^{-[n(n+\alpha+\beta+1)-m^2]t} P_n^{(\alpha, \beta)}(x) J_m(y).$$

In (4.7), putting $t = 0$, we have

$$(4.8) \quad f(x, y) = u(x, y, 0) = \sum_{\substack{n=0 \\ m=0}}^{\infty} A_{n,m} P_n^{(\alpha, \beta)}(x) J_m(y).$$

It is interesting to note that (2.2) is of the same form as (4.8).

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INDIRIZZO DELL'AUTORE:

S.D. Bajpai - Department of Mathematics - University of Bahrain - P.O. Box 32038, Isa Town - Bahrain